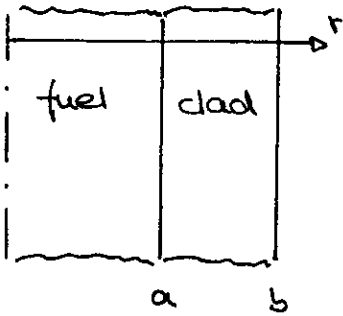


Solution - Problem Set No. 6

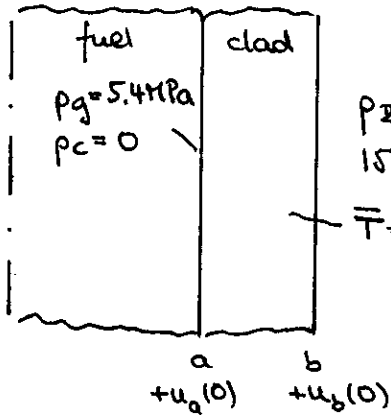
Geometry:



clad:
 inner radius: $a = 4.9 \text{ mm}$
 outer radius: $b = 5.6 \text{ mm}$
 wall thickness: $b - a = 0.7 \text{ mm}$

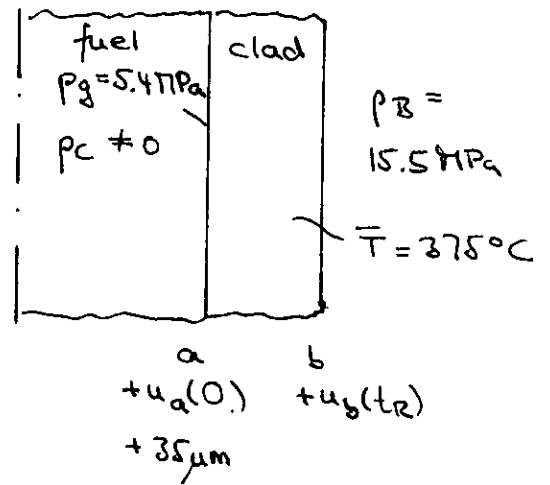
Operating Conditions:

hot zero power (time $t=0$)



no creep strains

hot full power (time $t = t_R$)



Computational Basis

- Finite-Element-treatment for clad: single ring
- $(u_a - u_a(0)) \propto t$ for each power ramp
- $(\bar{T} - \bar{T}_{H2P}) \propto t$ for each power ramp
- no axial force from fuel clad contact ($\bar{F}_{zc} = 0$)

From the class-notes, we get the following equations when treating the clad as a single ring:

$$\begin{pmatrix} \sigma_r \\ \sigma_\theta \\ \sigma_z \end{pmatrix} = \begin{pmatrix} \frac{-b}{a+b} & \frac{-a}{a+b} & 0 \\ -\frac{b}{b-a} & \frac{a}{b-a} & 0 \\ 0 & 0 & \frac{1}{b(b^2-a^2)} \end{pmatrix} \begin{pmatrix} p_b \\ p_a \\ \bar{T}_z \end{pmatrix} \quad (1)$$

with

$$p_b = p_3 \quad (2)$$

$$p_a = p_g + p_c \quad (3)$$

$$\bar{T}_z = \pi a^2 p_g - \pi b^2 p_3 \quad (4)$$

The strains and the displacements are related by

$$\begin{pmatrix} \epsilon_r \\ \epsilon_\theta \\ \epsilon_z \end{pmatrix} = \begin{pmatrix} \frac{1}{b-a} & -\frac{1}{b-a} & 0 \\ \frac{1}{a+b} & \frac{1}{a+b} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} u_b \\ u_a \\ \epsilon_z \end{pmatrix} \quad (5)$$

Otherwise, the strains can be calculated from the stresses, the thermal strain ϵ_T and the strains due to plasticity and creep $\epsilon_{mr}, \epsilon_{m\theta}, \epsilon_{mz}$.

$$\begin{pmatrix} \epsilon_r \\ \epsilon_\theta \\ \epsilon_z \end{pmatrix} = \underbrace{\frac{1}{E} \begin{pmatrix} 1 & -\nu & -\nu \\ -\nu & 1 & -\nu \\ -\nu & -\nu & 1 \end{pmatrix}}_{\text{elastic effects}} \begin{pmatrix} \sigma_r \\ \sigma_\theta \\ \sigma_z \end{pmatrix} + \underbrace{\epsilon_T \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}}_{\text{thermal}} + \underbrace{\begin{pmatrix} \epsilon_{mr} \\ \epsilon_{m\theta} \\ \epsilon_{mz} \end{pmatrix}}_{\text{plasticity and creep}} \quad (6)$$

Furthermore, for the strains due to creep we have the equations

$$\begin{pmatrix} \dot{\epsilon}_{mr} \\ \dot{\epsilon}_{m\theta} \\ \dot{\epsilon}_{mz} \end{pmatrix} = \frac{\dot{\epsilon}_g}{\sigma_g} \begin{pmatrix} 1 & -1/2 & -1/2 \\ -1/2 & 1 & -1/2 \\ -1/2 & -1/2 & 1 \end{pmatrix} \begin{pmatrix} \sigma_r \\ \sigma_\theta \\ \sigma_z \end{pmatrix} \quad (7)$$

with the equivalent von-Rises stress σ_g and

$$\dot{\epsilon}_g = 10^{-9} \frac{\sigma_g}{\text{MPa}}, \text{ so that eq. (7) is independent of } \sigma_g.$$

In order to use eq. (7) we have to derive :

- eq. (1)

$$\begin{pmatrix} \dot{q}_r \\ \dot{q}_\theta \\ \dot{q}_z \end{pmatrix} = \begin{pmatrix} \frac{-b}{a+b} & \frac{-a}{a+b} & 0 \\ \frac{-b}{b-a} & \frac{a}{b-a} & 0 \\ 0 & 0 & \frac{1}{r(b^2-a^2)} \end{pmatrix} \begin{pmatrix} \dot{p}_b \\ \dot{p}_a \\ \dot{f}_z \end{pmatrix}$$

and from eq. (2) to (4), with $p_a = \text{const.}$, $p_B = \text{const.}$

we get

$$\begin{aligned} \dot{p}_b &= 0 \\ \dot{p}_a &= \dot{p}_c \\ \dot{f}_z &= 0 \end{aligned}$$

and therefore

$$\begin{pmatrix} \dot{q}_r \\ \dot{q}_\theta \\ \dot{q}_z \end{pmatrix} = \begin{pmatrix} -\frac{a}{a+b} \\ \frac{a}{b-a} \\ 0 \end{pmatrix} \dot{p}_c \tag{8}$$

- eq. (5)

$$\begin{pmatrix} \dot{m}_r \\ \dot{m}_\theta \\ \dot{m}_z \end{pmatrix} = \begin{pmatrix} \frac{-1}{b-a} & \frac{-1}{b-a} & 0 \\ \frac{-1}{a+b} & \frac{1}{a+b} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \dot{u}_b \\ \dot{u}_a \\ \dot{E}_z \end{pmatrix} \tag{9}$$

- eq. (6)

$$\begin{pmatrix} \dot{E}_r \\ \dot{E}_\theta \\ \dot{E}_z \end{pmatrix} = \frac{1}{E} \begin{pmatrix} 1 & -v & -v \\ -v & 1 & -v \\ -v & -v & 1 \end{pmatrix} \begin{pmatrix} \dot{q}_r \\ \dot{q}_\theta \\ \dot{q}_z \end{pmatrix} + \dot{E}_T \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} + \begin{pmatrix} \dot{E}_{mr} \\ \dot{E}_{m\theta} \\ \dot{E}_{mz} \end{pmatrix} \tag{10}$$

Plugging eq. (7), (8) and (9) in eq. (10) yields:

(4)

$$\begin{pmatrix} \frac{1}{b-a} & \frac{-1}{b-a} & 0 \\ \frac{1}{a+b} & \frac{1}{a+b} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \ddot{u}_b \\ \ddot{u}_a \\ \dot{\epsilon}_z \end{pmatrix} = \frac{1}{E} \begin{pmatrix} 1 & -\nu & -\nu \\ -\nu & 1 & -\nu \\ -\nu & -\nu & 1 \end{pmatrix} \begin{pmatrix} \frac{-a}{a+b} \\ \frac{a}{b-a} \\ 0 \end{pmatrix} \dot{p}_c + \dot{\epsilon}_T \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \quad (11)$$

$$+ \beta \begin{pmatrix} 1 & -1/2 & -1/2 \\ -1/2 & 1 & -1/2 \\ -1/2 & -1/2 & 1 \end{pmatrix} \begin{pmatrix} \frac{-b}{a+b} & \frac{-a}{a+b} & 0 \\ \frac{-b}{b-a} & \frac{a}{b-a} & 0 \\ 0 & 0 & \frac{1}{\pi(b^2-a^2)} \end{pmatrix} \begin{pmatrix} p_B \\ p_B + p_c \\ \bar{F}_z \end{pmatrix}$$

$$\text{with } \beta = \frac{\dot{\epsilon}_T}{\sigma_g} = 10^{-9} \frac{1}{\text{MPa}} \quad (12)$$

writing eq. (11) in components gives:

$$\begin{aligned} \frac{1}{b-a} (\ddot{u}_b - \ddot{u}_a) &= -\frac{\dot{p}_c}{E} a \left(\frac{1}{a+b} + \frac{\nu}{b-a} \right) + \dot{\epsilon}_T + \frac{p_B}{2} \beta b \left(\frac{1}{b-a} - \frac{2}{a+b} \right) \\ &+ \frac{p_c + p_B}{2} \beta a \left(-\frac{1}{b-a} - \frac{2}{a+b} \right) - \beta \frac{\bar{F}_z}{2} \frac{1}{\pi(b^2-a^2)} \end{aligned} \quad (13)$$

$$\begin{aligned} \frac{1}{a+b} (\ddot{u}_b + \ddot{u}_a) &= \frac{\dot{p}_c}{E} a \left(\frac{\nu}{a+b} + \frac{1}{b-a} \right) + \dot{\epsilon}_T + \frac{p_B}{2} \beta b \left(\frac{1}{a+b} - \frac{2}{b-a} \right) \\ &+ \frac{p_c + p_B}{2} \beta a \left(\frac{1}{a+b} + \frac{2}{b-a} \right) - \beta \frac{\bar{F}_z}{2} \frac{1}{\pi(b^2-a^2)} \end{aligned} \quad (14)$$

$$\begin{aligned} \dot{\epsilon}_z &= \frac{\dot{p}_c}{E} a \left(\frac{\nu}{a+b} - \frac{\nu}{b-a} \right) + \dot{\epsilon}_T + \beta \frac{p_B}{2} b \left(\frac{1}{a+b} + \frac{1}{b-a} \right) \\ &+ \frac{p_c + p_B}{2} \beta a \left(\frac{1}{a+b} - \frac{1}{b-a} \right) + \beta \frac{\bar{F}_z}{2} \frac{1}{\pi(b^2-a^2)} \end{aligned} \quad (15)$$

Now, we have 3 equations (eq (13) to (15)) and 5 unknowns:

$$\ddot{u}_b, \ddot{u}_a, p_c \text{ and } \dot{p}_c, \dot{\epsilon}_T, \dot{\epsilon}_z$$

We need two more equations to calculate these unknowns.

(we only want to calculate $p_c(t)$, so we can reduce the problem to 2 equations (eq. (13) and (14)) with 4 unknowns

$$\ddot{u}_b, \ddot{u}_a, \dot{p}_c \text{ and } \dot{\epsilon}_T$$

From the calculational basis we know, that

(5)

$$\bullet [u_a - u_a(t=0)] \propto t$$

$$\text{and thereby } u_a(t) = u_a(t=0) + \frac{u_{aHFP} - u_a(t=0)}{t_R} t$$

$$\text{with } u_{aHFP} - u_{aHFP} = 35 \mu\text{m}.$$

$$\text{Therefore } \dot{u}_a = \frac{u_{aHFP} - u_{aHFP}}{t_R} \quad (16)$$

$$\bullet (\bar{T} - \bar{T}_{HFP}) \propto t$$

$$\text{and thereby with } \epsilon_T = \alpha (\bar{T}(t) - T_0) \quad \text{cold condition}$$

$$\epsilon_T(t) = \alpha [(\bar{T}_{HFP} - T_0) + \gamma t] \quad (17)$$

with

$$\gamma = \frac{\bar{T}_{HFP} - \bar{T}_{HFP}}{t_R} \quad (18)$$

and

$$\dot{\epsilon}_T = \alpha \gamma \quad (19)$$

Using eq. (16) and (19), eq. (13) and (14) reduce to two equations with two unknowns, \dot{u}_b and $p_c(t)$.

We can now eliminate \dot{u}_b from eq. (13) and (14) and get

$$2\dot{u}_a = \frac{1}{b^2 - a^2} \left[\frac{\dot{p}_c}{E} 2a (b^2(1+\nu) + a^2(1-\nu)) + 2a \dot{\epsilon}_T (b^2 - a^2) - p_B 4ab^2 + (p_c + p_g) \beta a (3b^2 + a^2) - \frac{\bar{T}_z}{\pi} \beta a \right] \quad (20)$$

Now, write eq. (20) in the form

$$\dot{p}_c + \varphi p_c = \psi \quad (21)$$

with

$$\zeta = \frac{E(b^2 - a^2)}{2a[b^2(1+\nu) + a^2(1-\nu)]}$$

$$\varphi = \frac{\beta a}{s^2 - a^2} (3b^2 + a^2) \zeta \quad (22)$$

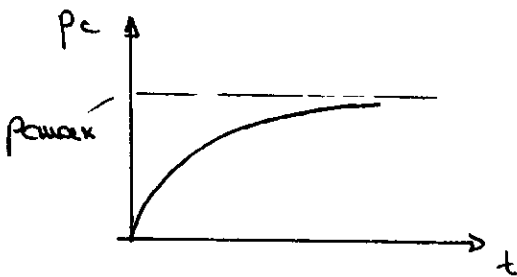
$$\psi = \zeta \left[\frac{1}{s^2 - a^2} \left(\frac{\beta F_z a}{\pi} + 4ab^2 \beta p_B - \beta a (3b^2 + a^2) p_g \right) + 2(\ddot{u}_a - a \ddot{\epsilon}_r) \right]$$

Equation (21) can be easily solved and has the solution

$$p_c(t) = c e^{-\varphi t} + \frac{\psi}{\varphi} \quad (23)$$

Applying the condition, that $p_c(t=0) = 0$ the constant c can be determined and eq. (23) becomes

$$p_c(t) = \frac{\psi}{\varphi} (1 - e^{-\varphi t}) \quad (24)$$



$$p_{cmax} = \frac{\psi}{\varphi}$$

Now, with $p_c(t=t_r) = p_{cr}$ from eq. (24) the clad stresses can be calculated using eq. (1)

$$\begin{pmatrix} \sigma_r \\ \sigma_\theta \\ \sigma_z \end{pmatrix}_{t=t_r} = \begin{pmatrix} -\frac{b}{a+b} & -\frac{a}{a+b} & 0 \\ -\frac{b}{b-a} & \frac{a}{b-a} & 0 \\ 0 & 0 & \frac{1}{\pi(b^2 - a^2)} \end{pmatrix} \begin{pmatrix} p_B \\ p_g + p_{cr} \\ \bar{F}_z \end{pmatrix} \quad (25)$$

In order to obtain the clad strains from eq. (6) we need to calculate the strains due to thermal and plasticity and Creep effects.

The thermal strains at hot full power are (also from eq. (17))

$$\epsilon_{T_R} = \epsilon_T (t = t_R) = \alpha (\bar{T}_{HFP} - \overset{\text{cold, shutdown}}{T_0}) \quad (26)$$

The plasticity and creep strains are determined by eq. (1) and (7):

$$\begin{pmatrix} \epsilon_{mr} \\ \epsilon_{m\theta} \\ \epsilon_{mz} \end{pmatrix} = \int_0^{t_R} \beta \begin{pmatrix} 1 & -1/2 & -1/2 \\ -1/2 & 1 & -1/2 \\ -1/2 & -1/2 & 1 \end{pmatrix} \begin{pmatrix} \frac{-b}{a+b} & \frac{-a}{a+b} & 0 \\ \frac{-b}{b-a} & \frac{a}{b-a} & 0 \\ 0 & 0 & \frac{1}{4(b^2-a^2)} \end{pmatrix} \begin{pmatrix} p_B \\ p_g + \frac{\psi}{\ell} (1 - e^{-\psi t}) \\ F_z \end{pmatrix} dt$$

$$= \frac{\beta}{2} \begin{pmatrix} -b^2 + 3ab & a^2 - 3ab & -\frac{1}{4} \\ -b^2 - 3ab & a^2 + 3ab & -\frac{1}{4} \\ 2b^2 & -2a^2 & +\frac{2}{\pi} \end{pmatrix} \frac{1}{b^2 - a^2} \begin{pmatrix} p_B t_R \\ p_g t_R + \frac{\psi}{\ell} (t_R - \frac{1}{\psi} (1 - e^{-\psi t_R})) \\ F_z t_R \end{pmatrix} \quad (27)$$

Applying eq. (25), (26) and (27) to eq. (6) yields the clad strains.

The radial displacement at the outer surface can be calculated by using eq. (5)

$$\begin{pmatrix} u_b \\ u_a \\ \epsilon_z \end{pmatrix}_{t=t_R} = \begin{pmatrix} \frac{1}{b-a} & -\frac{1}{b-a} & 0 \\ \frac{1}{a+b} & \frac{1}{a+b} & 0 \\ 0 & 0 & 1 \end{pmatrix}^{-1} \begin{pmatrix} \epsilon_r \\ \epsilon_\theta \\ \epsilon_z \end{pmatrix}_{t=t_R}$$

or, easier from

$$\left. \begin{aligned} \epsilon_r &= \frac{u_b - u_a}{b - a} \\ \epsilon_\theta &= \frac{u_b + u_a}{a + b} \end{aligned} \right\} u_b = \frac{1}{2} (\epsilon_r (b - a) + \epsilon_\theta (a + b)) \quad (28)$$

Results:

8

• $t_R = 30 \text{ min}$:

$$\begin{pmatrix} \sigma_r \\ \sigma_\theta \\ \sigma_z \end{pmatrix} = \begin{pmatrix} -38.9 \\ 336.0 \\ -48.5 \end{pmatrix} \text{ MPa}$$

$$\begin{pmatrix} \epsilon_r \\ \epsilon_\theta \\ \epsilon_z \end{pmatrix} = \begin{pmatrix} 0.8 \\ 7.4 \\ 0.6 \end{pmatrix} \times 10^{-3}$$

$$u_b = 3.91 \times 10^{-5} \text{ m}$$

• $t_R = 30 \text{ hours}$

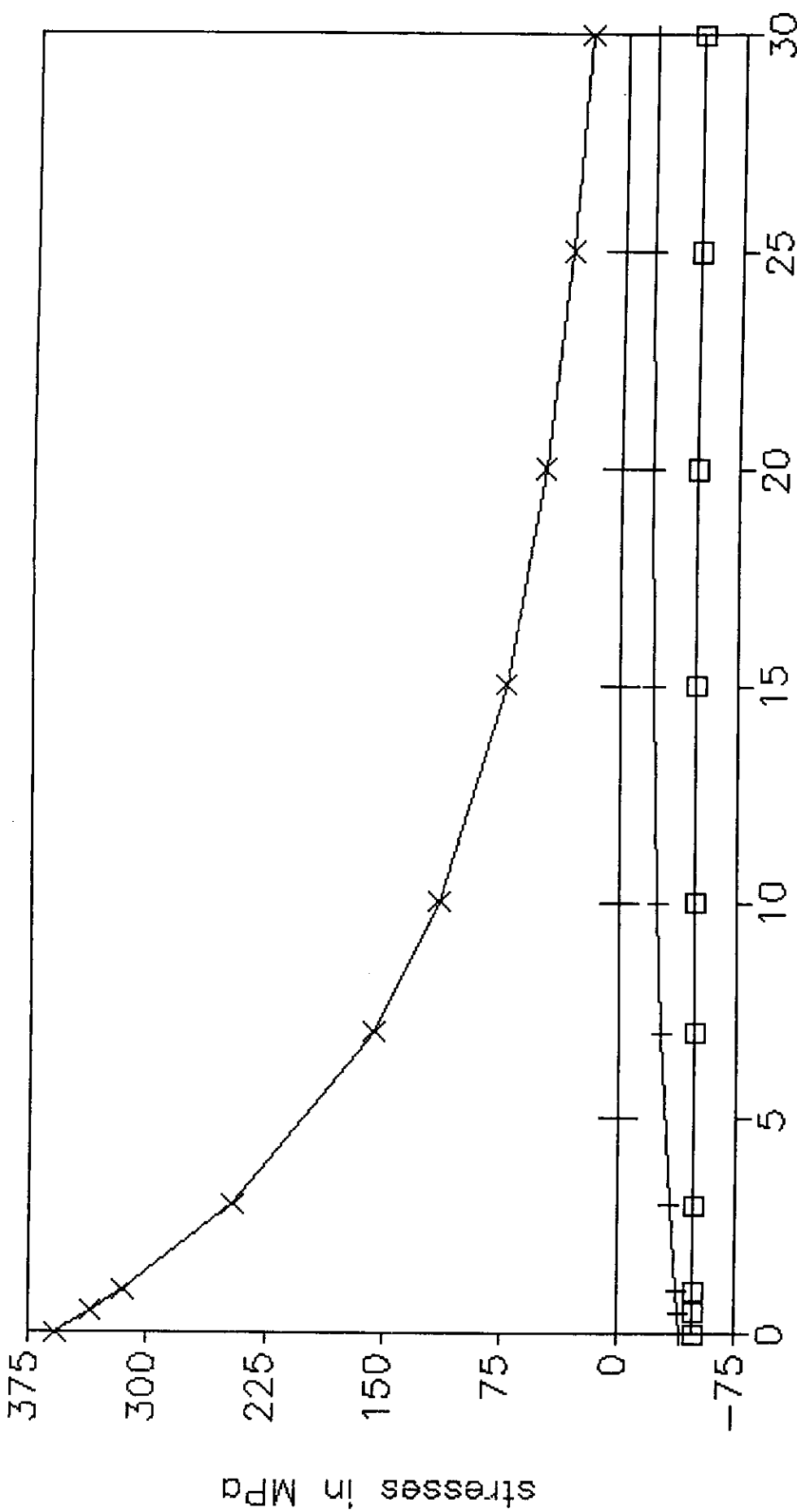
$$\begin{pmatrix} \sigma_r \\ \sigma_\theta \\ \sigma_z \end{pmatrix} = \begin{pmatrix} -18.0 \\ 22.6 \\ -48.5 \end{pmatrix} \text{ MPa}$$

$$\begin{pmatrix} \epsilon_r \\ \epsilon_\theta \\ \epsilon_z \end{pmatrix} = \begin{pmatrix} 2.5 \\ 7.5 \\ -4.8 \end{pmatrix} \times 10^{-3}$$

$$u_b = 4.02 \times 10^{-5} \text{ m}$$

Problem Set #6

Fuel-Clad Interaction stresses

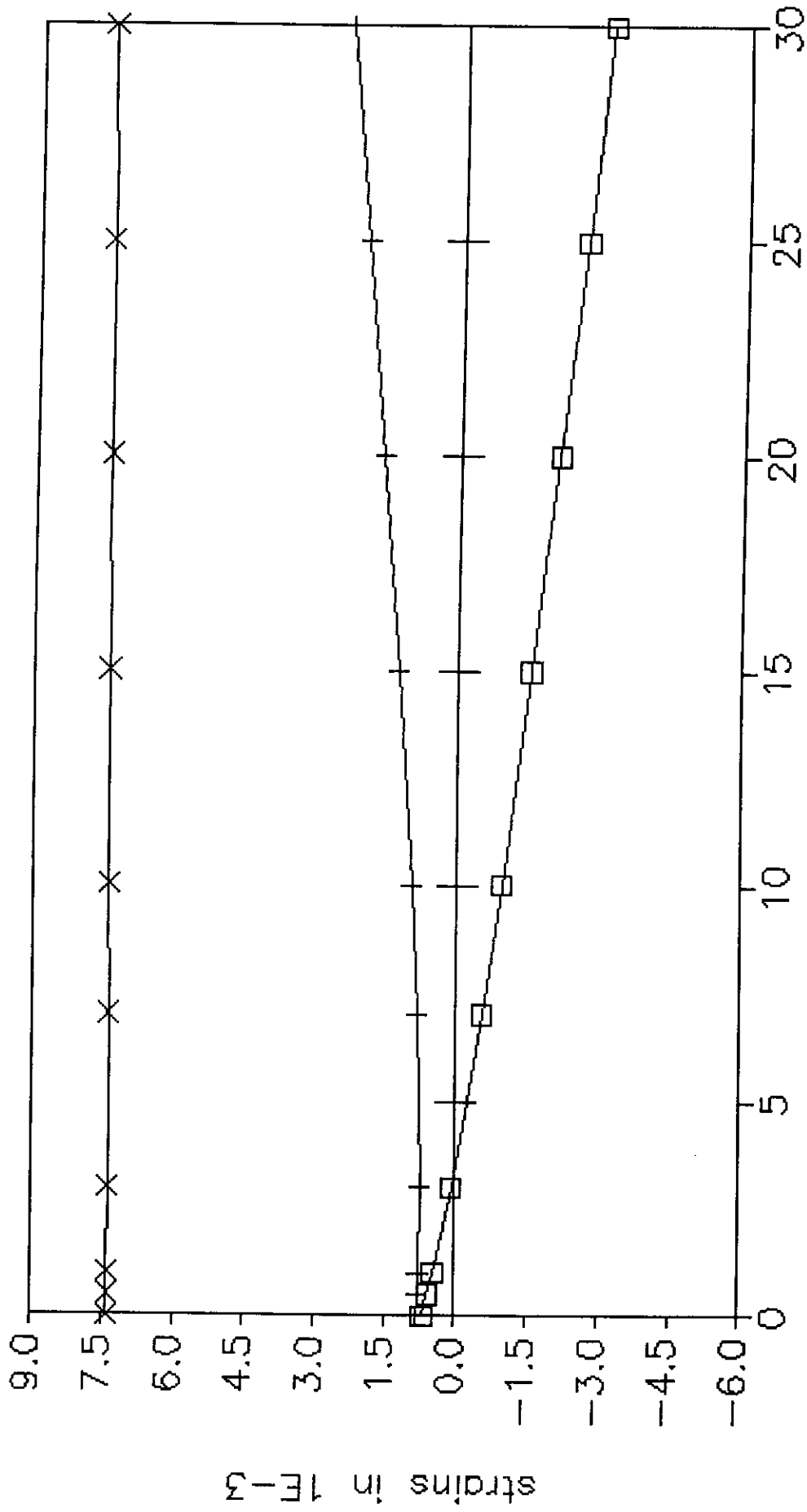


time in hours

+ stress in r-dir. x stress in t-dir. □ stress in z-dir.

creep strain rate:
eg = $1e-9$ * deviatoric stress

Problem Set #6
 Fuel-Clad Interaction
 strains



+ strain in r-dir. x strain in t-dir. o strain in z-dir.

creep strain rate:
 eg = $1e-9$ * deviatoric stress