

2.003J/1.053J Dynamics and Control I, Spring 2007
 Professor Thomas Peacock
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Lecture 3

Dynamics of a Single Particle: Angular Momentum

Example 2: Particle on String Pulled Through Hole

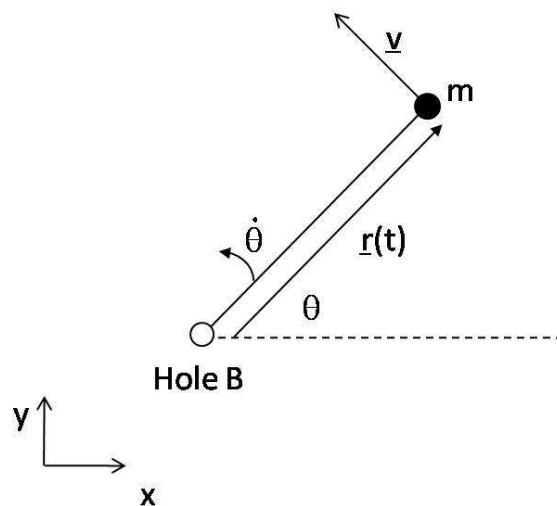


Figure 1: Particle on string pulled through hole. Tabletop with hole B. A string comes out with an attached mass. The particle is traveling around with an angular velocity $\dot{\theta}$. Figure by MIT OCW.

Assume: Frictionless surface. Inextensible String.

Pull string through hole at B such that:

$$\begin{aligned} r(t_0) &= L & \frac{dr}{dt}(t_0) &= 0 \\ r(t_1) &= L/2 & \frac{dr}{dt}(t_1) &= 0 \end{aligned}$$

If $\dot{\theta}(t_0) = \dot{\theta}_0$, what is $\dot{\theta}(t_1) = \dot{\theta}_1$?

Discussion

If we use linear momentum, will need to describe forces between m and string. Thinking about angular momentum about the point B:

$$\tau_B = \dot{\underline{h}}_B + \underline{v}_B \times m\underline{v} \leftarrow \text{Angular momentum principle}$$

$$\underline{h}_B = \underline{r} \times m\underline{v} = \text{Angular Momentum}$$

Now:

$$\tau_B = \underline{r} \times \underline{F} \leftarrow \text{Forces acting on particle } \tau_B = 0 \text{ because } \underline{r} \parallel \underline{F}$$

$$\tau_B = \dot{\underline{h}}_B + \underline{v}_B \times m\underline{v} \Rightarrow \text{Angular momentum about } B \text{ is constant } \dot{\underline{h}}_B = \underline{0}.$$

$$\tau_B = 0 \text{ (from above)}$$

$$\underline{v}_B = 0 \text{ because B is not moving}$$

$$\therefore h_B = \text{Constant}$$

In Cartesian Coordinates

$$\underline{r} = r \cos \theta \hat{i} + r \sin \theta \hat{j}$$

$$\underline{p} = m\underline{v} = m\dot{\underline{r}} = -mr\dot{\theta} \sin \theta \hat{i} + mr\dot{\theta} \cos \theta \hat{j}$$

a. $\underline{h}_B(t_0) = \underline{r} \times \underline{p} = LmL\dot{\theta}_0 \hat{k}$ (\hat{k} is unit vector in z-direction: out of page).

b. $\underline{h}_B(t_1) = \frac{L}{2} m \frac{L}{2} \dot{\theta}_1 \hat{k}$

Setting (a) = (b): $\dot{\theta}_1 = 4\dot{\theta}_0$, and velocity of particle $v_1 = 2v_0 = \frac{L}{2} 4\dot{\theta}_0 = 2L\dot{\theta}_0$.

Energy is not conserved: why? The pulling force (tension) does work.

Dynamics of systems of particlesForces on each particle may be composed as follows

$$\underline{F}_i = \underline{F}_i^{ext} + \underline{F}_i^{int}$$

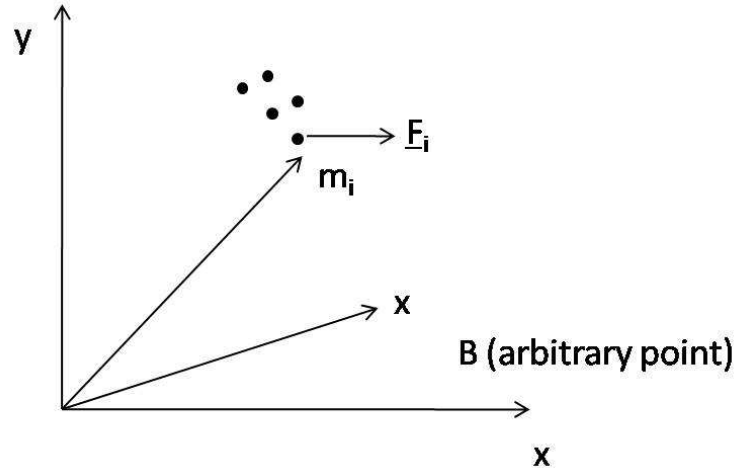


Figure 2: Dynamics of systems of particles. Figure by MIT OCW.

F_i : Resultant force acting on m_i

F_i^{ext} : External forces (e.g. gravity)

F_i^{int} : Internal forces between particles (e.g. charge attraction)

$$F_i^{int} = \sum_{j=1}^n \underline{f}_{ij} \text{ Force on particle } i \text{ due to particle } j$$

Newton's Third Law

$$\underline{f}_{ij} = -\underline{f}_{ji}$$

Thus:

$$\sum_{i=1}^n F_i^{int} = \sum_{i=1}^n \sum_{\substack{j=1 \\ j \neq i}}^n \underline{f}_{ij} = 0$$

Sum of all internal forces is zero, therefore:

$$\sum_{i=1}^n F_i^{int} = 0$$

Total internal torques is also zero: demonstrate by considering an arbitrary pair of particles:

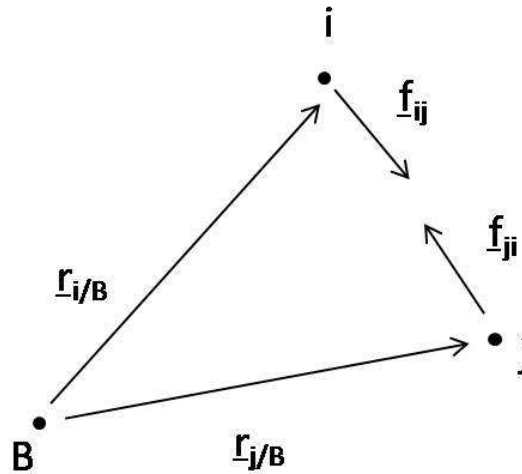


Figure 3: Arbitrary pair of particles subject to individual forces. Figure by MIT OCW.

$$\underline{\tau}_B = \underline{r}_{i/B} \times \underline{f}_{ij} + \underline{r}_{j/B} \times \underline{f}_{ji} = (\underline{r}_{i/B} - \underline{r}_{j/B}) \times \underline{f}_{ij}$$

$$\text{but } (\underline{r}_{i/B} - \underline{r}_{j/B} \parallel \underline{f}_{ij})$$

$$\therefore \underline{\tau}_B^{int} = 0 \text{ No net internal torque}$$

Center of mass

$$\underline{r}_c = \frac{\sum_{i=1}^n m_i \underline{r}_i}{\sum_{i=1}^n m_i} = \frac{\sum_{i=1}^n m_i \underline{r}_i}{M}$$

M : Total Mass of System

Note that this relation can also be written as $\sum_{i=1}^n m_i (\underline{r}_i - \underline{r}_c) = 0$ i.e. center of mass is the point about which the total mass moment is zero.

Newton's Laws for Systems of Particles

(Williams: C-1 to C-3.6)

Derivation needed to prevent mistakes in applying the laws later. Will be able to use results for rigid bodies.

Linear Momentum Principle (for a single particle)

$$\underline{F}_i = \frac{d}{dt} \underline{p}_i$$

\underline{F}_i : Total Force on particle i

\underline{p}_i : Linear momentum of particle i

$$\sum_{i=1}^n \underline{F}_i^{ext} + \sum_{i=1}^n \underline{F}_i^{int} = \frac{d}{dt} \sum_{i=1}^n \underline{p}_i = \frac{d}{dt} \underline{p}$$

$$\underline{F}_i^{int} = 0$$

$$\underline{F}^{ext} = \frac{d}{dt} \underline{p}$$

\underline{F}^{ext} : Sum of F external for whole system.

Note that total linear momentum:

$$\underline{p} = \sum_{i=1}^n \underline{p}_i = \sum_{i=1}^n m_i \underline{v}_i = M \underline{v}_c \text{ where } \underline{v}_c = \dot{\underline{r}}_c = \frac{d}{dt} \sum_{i=1}^n \frac{m_i \underline{r}_i}{M}$$

If $\sum_{i=1}^n \underline{F}_i^{ext} = \underline{F}^{ext} = 0 \Rightarrow \underline{p} = \text{constant}$; therefore, $\underline{v}_c = \text{constant}$.

Example: You have a ball as a ice skater. Throw object, both ball and skater move, but center of mass stays the same, does not move.

Angular Momentum Principle

From Newton II $\underline{F}_i = \frac{d}{dt} \underline{p}_i$

Torque:

$$\underline{r}'_i \times \underline{F}_i = \underline{r}'_i \times \frac{d}{dt} \underline{p}_i$$

Sum over all particles.

$$\sum_{i=1}^n \underline{r}'_i \times \underline{F}_i^{ext} = \sum_{i=1}^n \underline{r}'_i \times \frac{d}{dt} \underline{p}_i$$

Later will need vectors to center of mass.

$$\sum_{i=1}^n \underline{\tau}_{iB}^{ext} = \text{Sum of all external torques about B}$$

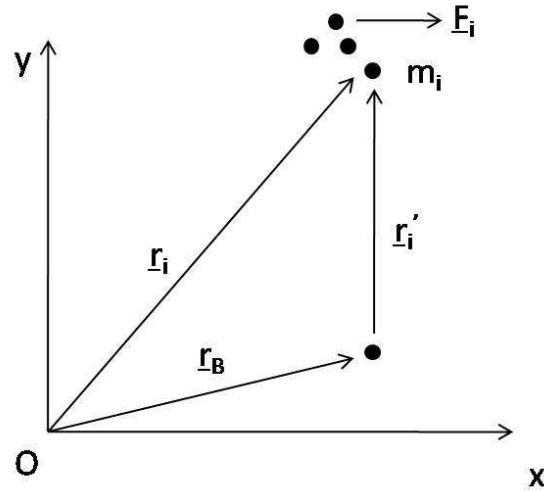


Figure 4: A system of particles subject to a force. Figure by MIT OCW.

$$\begin{aligned}\sum_{i=1}^n \mathcal{L}_{iB}^{ext} &= \sum_{i=1}^n \underline{r}'_i \times \frac{d}{dt} \underline{p}_i = \sum_{i=1}^n \frac{d}{dt} (\underline{r}'_i \times \underline{p}_i) - \sum_{i=1}^n \left(\frac{d}{dt} \underline{r}'_i \right) \times \underline{p}_i \\ \sum_{i=1}^n \mathcal{L}_{iB}^{ext} &= \frac{d}{dt} \sum_{i=1}^n \underline{h}_{iB} - \sum_{i=1}^n \frac{d}{dt} (\underline{r}_i - \underline{r}_B) \times \underline{p}_i \\ \tau_B^{ext} &= \frac{d}{dt} \underline{H}_B - \sum_{i=1}^n \underline{v}_i \times \underline{p}_i + \sum_{i=1}^n \underline{v}_B \times \underline{p}_i\end{aligned}$$

\underline{v}_B is the same for each \underline{p}_i .

$$\sum_{i=1}^n \underline{v}_B \times \underline{p}_i = \underline{v}_B \times \sum_{i=1}^n \underline{p}_i = \underline{v}_B \times \underline{P}$$

So, finally we have:

$$\mathcal{L}_B^{ext} = \frac{d}{dt} \underline{H}_B + \underline{v}_B \times \underline{P}$$

\mathcal{L}_B^{ext} : Total External Torque

$\frac{d}{dt} \underline{H}_B$: Total Angular Momentum

$\underline{v}_B \times \underline{P}$: Total Linear Momentum

Next time: Consequences of this expression and work-energy principle.